

Math 110A Midterm Solutions

C.T. Wildman

1. (a) The characteristic curves satisfy ODE

$$\frac{dy}{dx} = 2xe^y$$

Separating variables and integrating, we obtain

$$\int e^{-y} dy = \int 2x dx$$

Then $-e^{-y} = x^2 + C$ or $-C = x^2 + e^{-y}$. Thus, the general solution to our equation is $u(x, y) = f(x^2 + e^{-y})$ (for arbitrary differentiable f).

- (b) We have $u(x, 0) = f(x^2 + 1)$, and we want this to equal x^2 . Thus, we choose $f(s) = s - 1$, and our particular solution is then $u(x, y) = x^2 + e^{-y} - 1$.

- (c) Suppose $u(x, y) = h(x) + k(y)$ and plug in to the equation. This gives

$$h'(x) + 2xe^y k'(y) = x$$

Rearranging terms gives us

$$\frac{h'(x) - x}{2x} = -e^y k'(y)$$

However, the left hand side depends only on x whereas the right hand side depends only on y . Thus, both must be equal to some constant λ . This in turns gives us the pair of ODEs

$$\begin{aligned} h'(x) &= (2\lambda + 1)x \\ k'(y) &= -\lambda e^{-y} \end{aligned}$$

with particular solutions $h(x) = (\lambda + \frac{1}{2})x^2 = \lambda x^2 + \frac{1}{2}x^2$ and $k(y) = \lambda e^{-y}$. Thus, a particular solution to our equation is $\frac{1}{2}x^2 + \lambda(x^2 + e^{-y})$. We conclude that the general solution to the inhomogeneous equation is

$$u(x, y) = \frac{1}{2}x^2 + f(x^2 + e^{-y})$$

where f is any differentiable function (since the $\lambda(x^2 + e^{-y})$ term can be absorbed into f).

2. (a) We have $u(x, t) = \frac{1}{2} \int_{x-t}^{x+t} \psi(s) ds$, so it follows that $u(0, t) = \frac{1}{2} \int_{-t}^t \psi(s) ds$. Note that the value of the integral is the length of the intersection of $(-t, t)$ and $(-1, 1)$. So if $0 < t \leq 1$, we get $\frac{1}{2}(t - (-t)) = t$, but if $t > 1$, we get $\frac{1}{2}(1 - (-1)) = 1$. Summarizing, we have

$$u(0, t) = \begin{cases} t & 0 < t \leq 1 \\ 1 & t > 1 \end{cases}$$

- (b) Fix any $x \in (-\infty, \infty)$. Note that $u(x, t) = \frac{1}{2} \int_{x-t}^{x+t} \psi(s) ds$. The value of this integral is equal to the length of the intersection of $(x-t, x+t)$ and $(-1, 1)$. However, for t sufficiently large (to be exact: $t > \max\{x+1, -x+1\}$) the interval $(x-t, x+t)$ contains $(-1, 1)$ so we get

$$\lim_{t \rightarrow \infty} u(x, t) = \frac{1}{2}(1 - (-1)) = 1$$

- (c) We know that the energy $E(t)$ is conserved for the wave equation. In particular,

$$\lim_{t \rightarrow \infty} E(t) = E(0) = \frac{1}{2} \int_{-\infty}^{\infty} u_x(x, 0)^2 + u_t(x, 0)^2 dx = \frac{1}{2} \int_{-\infty}^{\infty} \psi(x)^2 dx = 1$$

3. (a)

$$u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-y^2} e^{-\frac{(x-y)^2}{4kt}} dy$$

- (b) Define

$$v(x, t) = \frac{1}{\sqrt{4kt+1}} e^{-\frac{x^2}{4kt+1}}$$

Taking derivatives shows that $v_t = kv_{xx}$ so $v(x, t)$ satisfies the heat equation. Furthermore, we have $v(x, 0) = e^{-x^2}$. Thus, by uniqueness, we have $v(x, t) = u(x, t)$ from (a) which is the desired result.

- (c) The easier case is the one with the Neumann condition $u_x(0, t) = 0$. To solve this, we would have used even reflection. However, our initial data is already even when considered as a function on the whole line. Therefore our solution formula remains unchanged and we get

$$u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-y^2} e^{-\frac{(x-y)^2}{4kt}} dy \quad \text{for } x \in [0, \infty)$$