# A New Recursion 

# In <br> The Theory of Macdonald Polynomials 

Joint work<br>with<br>Jim Haglund

## Some Basic Ingredients

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$$
n(\mu)=\sum l=\sum l^{\prime}=\sum(i-1) \mu_{i}=\sum\binom{\mu_{i}^{\prime}}{2}
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$\sum q^{a^{\prime}} t^{l^{\prime}}=B_{\mu}(q, t)$

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$n(\mu)=\sum l=\sum l^{\prime}=\sum(i-1) \mu_{i}=\sum\binom{\mu_{i}^{\prime}}{2}$
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$$
\Pi^{o, o}\left(1-q^{a^{\prime}} t^{l^{\prime}}\right)=\Pi_{\mu}(q, t)
$$

$$
\Pi\left(q^{a}-t^{l+1}\right)=\tilde{h}_{\mu}(q, t)
$$

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$$
\begin{aligned}
& n(\mu)=\sum l=\sum l^{\prime}=\sum(i-1) \mu_{i}=\sum\left(\begin{array}{c}
\left.\mu_{2}^{\prime}\right) \\
\sum q^{a^{\prime}} t^{l^{\prime}}
\end{array}=B_{\mu}(q, t) \quad \prod^{o, o}\left(1-q^{a^{\prime}} t^{l^{\prime}}\right)=\Pi_{\mu}(q, t)\right.
\end{aligned}
$$

$$
\Pi\left(q^{a}-t^{l+1}\right)=\tilde{h}_{\mu}(q, t)
$$

$$
\Pi\left(t^{l}-q^{a+1}\right)=\tilde{h}_{\mu}^{\prime}(q, t)
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\begin{aligned}
n(\mu)=\sum l & =\sum l^{\prime}=\sum(i-1) \mu_{i}=\sum\left(\begin{array}{c}
\mu_{2}^{\prime} i
\end{array}\right) \\
\sum q^{a^{\prime}} t^{l^{\prime}} & =B_{\mu}(q, t) \quad \Pi^{o, o}\left(1-q^{a^{\prime}} t^{l^{\prime}}\right)=\Pi_{\mu}(q, t)
\end{aligned}
$$

$$
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& \Pi\left(q^{a}-t^{l+1}\right)=\tilde{h}_{\mu}(q, t) \\
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M=(1-t)(1-q)
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The Macdonald polynomials

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the unique symmetric function basis $\left\{\tilde{\mathrm{H}}_{\mu}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]\right\}_{\mu}$

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$$
\text { 1) }\left.\quad \tilde{\mathrm{H}}_{\mu}[\mathbf{X} ; \boldsymbol{q}, \mathrm{t}]\right|_{\mathbb{S}_{[n]}}=1
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2) $\left\langle\tilde{\mathrm{H}}_{\mu}, \tilde{\mathrm{H}}_{\lambda}\right\rangle_{*}=\chi(\lambda=\mu) \mathbf{w}_{\mu}(\mathbf{q}, \mathbf{t})$

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\left\langle p_{\alpha}, p_{\beta}\right\rangle_{*}= \\
\chi(\alpha=\beta) z_{\alpha} \prod_{\alpha_{i}>0}(-1)^{\alpha_{i}-1}\left(1-t^{\alpha_{i}}\right)\left(1-t^{\beta_{i}}\right)
\end{gathered}
$$

The Visual Representation of a polynomial P[q, t]

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\text { The Visual Representation of a polynomial } \mathrm{P}[\mathrm{q}, \mathrm{t}] \\
8 q^{3}+36 q^{2} t+54 q t^{2}+27 t^{3}+12 q^{2}+36 q t+27 t^{2}+6 q+9 t+1
\end{gathered}
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$$
8 q^{3}+36 q^{2} t+54 q t^{2}+27 t^{3}+12 q^{2}+36 q t+27 t^{2}+6 q+9 t+1
$$

$$
\left[\begin{array}{rrrr}
27 & 0 & 0 & 0 \\
27 & 54 & 0 & 0 \\
9 & 36 & 36 & 0 \\
1 & 6 & 12 & 8
\end{array}\right]
$$

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$$

$$
\mathfrak{t}^{\mathbf{r}} \begin{array}{r}
3 \\
2
\end{array}\left[\begin{array}{rrrr}
27 & 0 & 0 & 0 \\
1 \\
27 & 54 & 0 & 0 \\
0 & 36 & 36 & 0 \\
1 & 6 & 12 & 8
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The Visual Representation of a polynomial P[q, t]

$$
\begin{gathered}
8 q^{3}+36 q^{2} t+54 q t^{2}+\sqrt{27} t^{3}+12 q^{2}+36 q t+27 t^{2}+6 q+9 t+1 \\
\left.\mathrm{t}^{\mathrm{r}}\left|\begin{array}{c}
3 \\
2 \\
1 \\
0
\end{array}\right| \begin{array}{rrrrr}
27 & 0 & 0 & 0 \\
27 & 54 & 0 & 0 \\
9 & 36 & 36 & 0 \\
1 & 6 & 12 & 8
\end{array}\right] \\
\xrightarrow[0]{2} 1 \\
\mathbf{q}^{\mathrm{S}}
\end{gathered}
$$

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The Visual Representation of a polynomial P[q, t]

Our Macdonald polynomials have Schur function expansion

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$$
\tilde{H}_{\mu}(x ; q, t)=\sum_{\lambda} S_{\lambda}(x) \tilde{K}_{\lambda, \mu}(q, t)
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$$

where

$$
\tilde{\mathbf{K}}_{\lambda \mu}(\boldsymbol{q}, \mathrm{t})=\mathrm{t}^{\mathbf{n}(\mu)} \mathbf{K}_{\lambda \mu}(\mathbf{q}, 1 / \mathbf{t})
$$

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$$

where

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\tilde{\mathbf{K}}_{\lambda \mu}(\boldsymbol{q}, \mathrm{t})=\mathrm{t}^{\mathbf{n}(\mu)} \mathrm{K}_{\lambda \mu}(\boldsymbol{q}, 1 / \mathrm{t})
$$

with $K_{\lambda_{\mu}}(\mathbf{q}, \mathrm{t})$ the Macdonald q , t -Kostka coefficient.

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$$
\tilde{H}_{[3,2]}(x ; q, t)=
$$

Our Macdonald polynomials have Schur function expansion

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with $K_{\lambda_{\mu}}(\mathbf{q}, \mathrm{t})$ the Macdonald $\mathrm{q}, \mathrm{t}$-Kostka coefficient.

$$
\tilde{H}_{[3,2]}(x ; q, t)=s_{5}+s_{4,1}\left[\begin{array}{lll}
1 & 1 & 0 \\
0 & 1 & 1
\end{array}\right]+s_{3,2}\left[\begin{array}{llll}
1 & 0 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 1 & 1
\end{array}\right]+s_{3,1,1}\left[\begin{array}{lll}
1 & 0 & 0 \\
1 & 2 & 1 \\
0 & 0 & 1
\end{array}\right]+
$$

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with $\mathrm{K}_{\lambda_{\mu}}(\mathrm{q}, \mathrm{t})$ the Macdonald q , t-Kostka coefficient.

$$
\begin{aligned}
& \tilde{H}_{[3,2]}(x ; q, t)=s_{5}+s_{4,1}\left[\begin{array}{lll}
1 & 1 & 0 \\
0 & 1 & 1
\end{array}\right]+s_{3,2}\left[\begin{array}{cccc}
1 & 0 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 1 & 1
\end{array}\right]+s_{3,1,1 q}\left[\begin{array}{lll}
1 & 0 & 0 \\
1 & 2 & 1 \\
0 & 0 & 1
\end{array}\right]+ \\
& +s_{2,2,1} q\left[\begin{array}{llll}
1 & 1 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 0 & 1
\end{array}\right]+s_{2,1,1,1} t q^{2}\left[\begin{array}{lll}
1 & 1 & 0 \\
0 & 1 & 1
\end{array}\right]+s_{1,1,1,1,1} \quad t^{2} q^{4}
\end{aligned}
$$

Our Macdonald polynomials have Schur function expansion

$$
\tilde{H}_{\mu}(x ; q, t)=\sum_{\lambda} S_{\lambda}(x) \tilde{K}_{\lambda, \mu}(q, t)
$$

where

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\tilde{\mathbf{K}}_{\lambda \mu}(\boldsymbol{q}, \mathrm{t})=\mathrm{t}^{\mathbf{n}(\mu)} \mathbf{K}_{\lambda \mu}(\mathbf{q}, 1 / \mathbf{t})
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$$
\begin{array}{r}
\tilde{H}_{[3,2]}(x ; q, t)=s_{5}+s_{4,1}\left[\begin{array}{lll}
1 & 1 & 0 \\
0 & 1 & 1
\end{array}\right]+s_{3,2}\left[\begin{array}{llll}
1 & 0 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 1 & 1
\end{array}\right]+s_{3,1,1}\left[\begin{array}{lll}
1 & 0 & 0 \\
1 & 2 & 1 \\
0 & 0 & 1
\end{array}\right]+ \\
+s_{2,2,1}
\end{array} \begin{aligned}
& q\left[\begin{array}{llll}
1 & 1 & 0 & 0 \\
0 & 1 & 1 & 0 \\
0 & 0 & 0 & 1
\end{array}\right]+s_{2,1,1,1} t q^{2}\left[\begin{array}{lll}
1 & 1 & 0 \\
0 & 1 & 1
\end{array}\right]+s_{1,1,1,1,1}
\end{aligned} t^{2} q^{4},
$$

## The Jig Saw Puzzle

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of

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of
Schur Function Expansions

## The Jig Saw Puzzle

$\tilde{\mathrm{H}}_{[\mathbf{1 , 1 , 1 , 1 ]}}[\mathbf{X} ; \boldsymbol{q}, \mathrm{t}]$
of
Schur Function Expansions

## The Jig Saw Puzzle

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{1 , 1 , 1 , 1}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{c}
s_{1,1,1,1} \\
s_{2,1,1} \\
s_{2,2}+s_{2,1,1} \\
s_{3,1}+s_{2,1,1} \\
s_{3,1}+s_{2,2} \\
s_{3,1} \\
s_{4}
\end{array}\right]}
\end{gathered}
$$

of

Schur Function Expansions

## The Jig Saw Puzzle



## The Jig Saw Puzzle

$\tilde{\mathrm{H}}_{[\mathbf{1 , 1 , 1 , 1 ]}}[\mathbf{X} ; \mathbf{q}, \mathbf{t}]$
\(\left[\begin{array}{c}s_{1,1,1,1} <br>

s_{2,1,1}\end{array}\right]\)| Frobenius characteristic of |
| :---: |
| the linear span |
| of derivatives of |

the Vandermonde determinant

$$
\tilde{\mathrm{H}}_{[\mathbf{2}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]
$$

## The Jig Saw Puzzle

$$
\begin{aligned}
& \tilde{\mathrm{H}}_{[\mathbf{1}, \mathbf{1}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& {\left[\begin{array}{c}
s_{1,1,1}, 1 \\
s_{2,1}, 1 \\
s_{2,2}+s_{2,1,1} \\
s_{3,1}+s_{2,1,1} \\
s_{3,1}+s_{2,2} \\
s_{3,1} \\
s_{4}
\end{array}\right]}
\end{aligned}
$$

Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant
of

## Schur Function Expansions

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

## The Jig Saw Puzzle

$$
\begin{aligned}
& \tilde{\mathrm{H}}_{[\mathbf{1 , 1 , 1 , 1} \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& {\left[\begin{array}{c}
s_{1,1,1,1} \\
s_{2,1,1} \\
s_{2,2}+s_{2,1,1} \\
s_{3,1}+s_{2,1,1} \\
s_{3,1}+s_{2,2} \\
s_{3,1} \\
s_{4}
\end{array}\right]}
\end{aligned}
$$

Frobenius characteristic of the linear span of derivatives of
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of
Schur Function Expansions

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]} \\
\\
\quad \tilde{\mathrm{H}}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}]
\end{gathered}
$$

## The Jig Saw Puzzle

$$
\begin{aligned}
& \tilde{\mathrm{H}}_{[\mathbf{1}, \mathbf{1}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& {\left[\begin{array}{c}
s_{1,1}, 1,1 \\
s_{2,1,1} \\
s_{2,2}+s_{2,1,1} \\
s_{3,1}+s_{2,1,1} \\
s_{3,1}+s_{2,2} \\
s_{3,1} \\
s_{4}
\end{array}\right]}
\end{aligned}
$$

Frobenius characteristic of the linear span of derivatives of
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## Schur Function Expansions

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 1 , 1}]}\left[\mathbf{X} ; \mathbf{\mathbf { q } _ { 1 }}, \mathrm{t}\right] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]} \\
\tilde{H}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{c}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]}
\end{gathered}
$$

## The Jig Saw Puzzle



Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant
of

## Schur Function Expansions

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

$$
\begin{gathered}
c \tilde{H}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1} 1 & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]}
\end{gathered}
$$

$$
\tilde{\mathrm{H}}_{[3,1]}[\mathbf{X} ; \mathbf{c}, \mathrm{t}]
$$

## The Jig Saw Puzzle



Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant
of

## Schur Function Expansions

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 1 , 1 ]}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]} \\
\tilde{\mathrm{H}}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]} \\
{\left[\begin{array}{cccc}
s_{3,1} & s_{2,2}+s_{2,1,1} & s_{2,1,1} & s_{1,1,1,1} \\
s_{4} & s_{3,1} & s_{3,1}+s_{2,2} & s_{2,1,1}
\end{array}\right]}
\end{gathered}
$$

## The Jig Saw Puzzle



Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant
of

## Schur Function Expansions

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 1 , \mathbf { 1 } ]}}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

$$
\begin{gathered}
c \tilde{\mathrm{H}}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]}
\end{gathered}
$$

$$
\tilde{\mathrm{H}}_{[3,1]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]
$$

$$
\left[\begin{array}{cccc}
s_{3,1} & s_{2,2}+s_{2,1,1} & s_{2,1,1} & s_{1,1,1,1} \\
s_{4} & s_{3,1} & s_{3,1}+s_{2,2} & s_{2,1,1}
\end{array}\right]
$$

$$
\tilde{\mathrm{H}}_{[2,2]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]
$$

## The Jig Saw Puzzle

$$
\begin{aligned}
& \tilde{H}_{[\mathbf{1}, \mathbf{1}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& {\left[\begin{array}{c}
s_{1,1}, 1,1 \\
s_{2,1} \\
s_{2,2}+s_{2,1,1} \\
s_{3,1}+s_{2,1,1} \\
s_{3,1}+s_{2,2} \\
s_{3,1} \\
s_{4}
\end{array}\right]}
\end{aligned}
$$

Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant
of

## Schur Function Expansions

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2}, \mathbf{1}, \mathbf{1}]}\left[\mathbf{X} ; \mathbf{c}_{\mathbf{\prime}} \mathbf{t}\right] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

$$
\begin{gathered}
c \tilde{\mathrm{H}}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[s_{2,2}\right.} \\
s_{2,1,1}
\end{gathered} s_{1,1,1,1}\left[\begin{array}{ccc}
s_{3,1} & s_{3,1}+s_{2,1} 1 & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]
$$

$$
\tilde{\mathrm{H}}_{[3, \mathbf{1}]}[\mathbf{X} ; \mathbf{c}, \mathrm{t}]
$$

$$
\left[\begin{array}{cccc}
s_{3,1} & s_{2,2}+s_{2,1,1} & s_{2,1,1} & s_{1,1,1,1} \\
s_{4} & s_{3,1} & s_{3,1}+s_{2,2} & s_{2,1,1}
\end{array}\right]
$$

$$
\begin{gathered}
\tilde{H}_{[\mathbf{2 , 2}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[s_{4}, s_{3,1}, s_{3,1}+s_{2,2}, s_{3,1}+s_{2,1,1}, s_{2,2}+s_{2,1,1}, s_{2,1,1}, s_{1,1,1,1}\right]}
\end{gathered}
$$

## The Jig Saw Puzzle

$$
\begin{aligned}
& \tilde{\mathrm{H}}_{[\mathbf{1 , 1 , 1 , \mathbf { 1 } ]}}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& {\left[\begin{array}{c}
s_{1,1,1,1} \\
s_{2,1,1} \\
s_{2,2}+s_{2,1,1} \\
s_{3,1}+s_{2,1,1} \\
s_{3,1}+s_{2,2} \\
s_{3,1} \\
s_{4}
\end{array}\right]}
\end{aligned}
$$

Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant
of

## Schur Function Expansions

They are all deformations
the Sn Harmonics

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[4]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]} \\
\\
\end{gathered}
$$

$$
\tilde{\mathrm{H}}_{[2,2]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]
$$

$$
\left[s_{4}, s_{3,1}, s_{3,1}+s_{2,2}, s_{3,1}+s_{2,1,1}, s_{2,2}+s_{2,1,1}, s_{2,1,1}, s_{1,1,1,1}\right]
$$

## The Jig Saw Puzzle

$$
\begin{aligned}
& \tilde{\mathrm{H}}_{[1,1,1,1]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& {\left[\begin{array}{c}
s_{1,1,1,1} \\
s_{2,1,1} \\
s_{2,2}+s_{2,1,1} \\
s_{3,1}+s_{2,1,1} \\
s_{3,1}+s_{2,2} \\
s_{3,1} \\
s_{4}
\end{array}\right]}
\end{aligned}
$$

Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant

# of <br> Schur Function Expansions 

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 1 , 1}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

What makes<br>the Schur functions<br>move???

They are all

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[4]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]} \\
\\
\end{gathered}
$$ deformations

the Sn Harmonics

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 2}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[s_{4}, s_{3,1}, s_{3,1}+s_{2,2}, s_{3,1}+s_{2,1,1}, s_{2,2}+s_{2,1,1}, s_{2,1,1}, s_{1,1,1,1}\right]}
\end{gathered}
$$

## The Jig Saw Puzzle



Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant

# of <br> Schur Function Expansions 

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 1 , \mathbf { 1 } ]}}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

What makes<br>the Schur functions<br>move???

Macdonald Reciprocity

They are all

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[4]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]} \\
\\
\end{gathered}
$$ deformations

the Sn Harmonics

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 2}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[s_{4}, s_{3,1}, s_{3,1}+s_{2,2}, s_{3,1}+s_{2,1,1}, s_{2,2}+s_{2,1,1}, s_{2,1,1}, s_{1,1,1,1}\right]}
\end{gathered}
$$

## The Jig Saw Puzzle

$\tilde{\mathrm{H}}_{[\mathbf{1 , 1 , 1 , \mathbf { 1 } ]}}[\mathbf{X} ; \mathbf{q}, \mathbf{t}]$
$\left[\begin{array}{c}s_{1,1,1,1} \\ s_{2,1,1} \\ s_{2,2}+s_{2,1,1} \\ s_{3,1}+s_{2,1,1} \\ s_{3,1}+s_{2,2} \\ s_{3,1} \\ s_{4}\end{array}\right]$

Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 1 , 1}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

of

## Schur Function Expansions

What makes<br>the Schur functions move???

## Macdonald Reciprocity

They are all deformations
the Sn Harmonics

$$
\tilde{\mathrm{H}}_{[2,2]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]
$$

$$
\left[s_{4}, s_{3,1}, s_{3,1}+s_{2,2}, s_{3,1}+s_{2,1,1}, s_{2,2}+s_{2,1,1}, s_{2,1,1}, s_{1,1,1,1}\right]
$$

$$
\begin{aligned}
& \frac{\tilde{\mathrm{H}}_{\mu}\left[\mathrm{B}_{\lambda}(\mathbf{q}, \mathrm{t})\right]}{\Pi_{\mu}(\mathbf{q}, \mathrm{t})}=\frac{\tilde{\mathrm{H}}_{\lambda}\left[\mathrm{B}_{\mu}(\mathbf{q}, \mathrm{t})\right]}{\Pi_{\lambda}(\mathbf{q}, \mathrm{t})} \\
& \tilde{H}_{[4]}[\mathbf{X} ; q, t] \\
& \begin{array}{c}
c \mathrm{H}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]}
\end{array} \\
& \tilde{H}_{[3,1]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& {\left[\begin{array}{cccc}
s_{3,1} & s_{2,2}+s_{2,1,1} & s_{2,1,1} & s_{1,1,1,1} \\
s_{4} & s_{3,1} & s_{3,1}+s_{2,2} & s_{2,1,1}
\end{array}\right]}
\end{aligned}
$$

## The Jig Saw Puzzle

$\tilde{\mathrm{H}}_{[\mathbf{1}, \mathbf{1}, \mathbf{1}, \mathbf{1}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}]$
$\left[\begin{array}{c}s_{1,1,1,1} \\ s_{2,1,1} \\ s_{2,2}+s_{2,1,1} \\ s_{3,1}+s_{2,1,1} \\ s_{3,1}+s_{2,2} \\ s_{3,1} \\ s_{4}\end{array}\right]$

Frobenius characteristic of the linear span of derivatives of
the Vandermonde determinant

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[\mathbf{2 , 1 , 1}]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[\begin{array}{cc}
s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1}+s_{2,2} & s_{2,1,1} \\
s_{3,1} & s_{2,2}+s_{2,1,1} \\
s_{4} & s_{3,1}
\end{array}\right]}
\end{gathered}
$$

of

## Schur Function Expansions

What makes<br>the Schur functions move???

## Macdonald Reciprocity

$$
\frac{\tilde{\mathrm{H}}_{\mu}\left[\mathrm{B}_{\lambda}(\mathbf{q}, \mathrm{t})\right]}{\Pi_{\mu}(\mathbf{q}, \mathrm{t})}=\frac{\tilde{\mathrm{H}}_{\lambda}\left[\mathrm{B}_{\mu}(\mathbf{q}, \mathrm{t})\right]}{\Pi_{\lambda}(\mathbf{q}, \mathrm{t})}
$$

They are all deformations
the Sn Harmonics

$$
\begin{gathered}
c \tilde{\mathrm{H}}_{[\mathbf{4}]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{ccc}
s_{2,2} & s_{2,1,1} & s_{1,1,1,1} \\
s_{3,1} & s_{3,1}+s_{2,1,1} & s_{2,1,1} \\
s_{4} & s_{3,1} & s_{2,2}
\end{array}\right]}
\end{gathered}
$$

## What else ?

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[3,1]}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
{\left[\begin{array}{cccc}
s_{3,1} & s_{2,2}+s_{2,1,1} & s_{2,1,1} & s_{1,1,1,1} \\
s_{4} & s_{3,1} & s_{3,1}+s_{2,2} & s_{2,1,1}
\end{array}\right]}
\end{gathered}
$$

$$
\begin{gathered}
\tilde{\mathrm{H}}_{[2,2]}[\mathbf{X} ; \mathbf{q}, \mathbf{t}] \\
{\left[s_{4}, s_{3,1}, s_{3,1}+s_{2,2}, s_{3,1}+s_{2,1,1}, s_{2,2}+s_{2,1,1}, s_{2,1,1}, s_{1,1,1,1}\right]}
\end{gathered}
$$

## k-Schur expansion of Macdonald Polynomials

## k-Schur expansion of Macdonald Polynomials

(for 2 bounded partitions of 4 )

## k-Schur expansion of Macdonald Polynomials

(for 2 bounded partitions of 4 )

H([2,2])

## k-Schur expansion of Macdonald Polynomials

(for 2 bounded partitions of 4 )
$\mathrm{H}([2,2]) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2}$

## k-Schur expansion of Macdonald Polynomials

(for 2 bounded partitions of 4 )

$$
\mathrm{H}([2,2]) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right]
$$

## k-Schur expansion of Macdonald Polynomials

 (for 2 bounded partitions of 4 )$$
\begin{gathered}
\mathrm{H}([2,2]) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right] \\
H([2,1,1]
\end{gathered}
$$

## k-Schur expansion of Macdonald Polynomials

 (for 2 bounded partitions of 4 )$$
\begin{array}{r}
\mathrm{H}([2,2]) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right] \\
H\left([2,1,1] \quad t q A_{1,1,1,1}+\left(q+t^{2}\right) A_{2,1,1}+A_{2,2}\right.
\end{array}
$$

## k-Schur expansion of Macdonald Polynomials

 (for 2 bounded partitions of 4 )$$
\begin{aligned}
& \mathrm{H}\left([\mathbf{2 , 2 ]}) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right]\right. \\
& \mathrm{H}\left([2,1,1] \quad t q A_{1,1,1,1}+\left(q+t^{2}\right) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{cc}
A_{2,1,1} & 0 \\
0 & A_{1,1,1,1} \\
A_{2,2} & A_{2,1,1}
\end{array}\right]\right.
\end{aligned}
$$

## k-Schur expansion of Macdonald Polynomials

 (for 2 bounded partitions of 4 )$$
\left.\begin{array}{cc}
\mathrm{H}([2,2]) & q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2}
\end{array} \begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right]\left(\begin{array}{cc} 
\\
H\left([2,1,1] \quad t q A_{1,1,1,1}+\left(q+t^{2}\right) A_{2,1,1}+A_{2,2}\right.
\end{array}\left[\begin{array}{cc}
A_{2,1,1} & 0 \\
0 & A_{1,1,1,1} \\
A_{2,2} & A_{2,1,1}
\end{array}\right]\right)
$$

## k-Schur expansion of Macdonald Polynomials

 (for 2 bounded partitions of 4 )$$
\left.\left.\begin{array}{l}
\mathrm{H}([2,2]) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right] \\
H\left([2,1,1] \quad t q A_{1,1,1,1}+\left(q+t^{2}\right) A_{2,1,1}+A_{2,2}\right.
\end{array} \begin{array}{ccc}
A_{2,1,1} & 0 \\
0 & A_{1,1,1,1} \\
A_{2,2} & A_{2,1,1}
\end{array}\right]\right)
$$

## k-Schur expansion of Macdonald Polynomials

 (for 2 bounded partitions of 4 )$$
\begin{aligned}
& \mathrm{H}\left([\mathbf{2 , 2 ]}) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right]\right. \\
& \mathrm{H}\left([2,1,1] \quad t q A_{1,1,1,1}+\left(q+t^{2}\right) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{cc}
A_{2,1,1} & 0 \\
0 & A_{1,1,1,1} \\
A_{2,2} & A_{2,1,1}
\end{array}\right]\right. \\
& \mathrm{H}([1,1,1,1]) \quad t^{4} A_{1,1,1,1}+\left(t^{2}+t^{3}\right) A_{2,1,1}+A_{2,2}
\end{aligned}
$$

## k-Schur expansion of Macdonald Polynomials

 (for 2 bounded partitions of 4 )$$
\begin{aligned}
& \mathrm{H}\left([\mathbf{2 , 2 ]}) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right]\right. \\
& \mathrm{H}\left([2,1,1] \quad t q A_{1,1,1,1}+\left(q+t^{2}\right) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{cc}
A_{2,1,1} & 0 \\
0 & A_{1,1,1,1} \\
A_{2,2} & A_{2,1,1}
\end{array}\right]\right. \\
& \mathrm{H}([1,1,1,1]) \quad t^{4} A_{1,1,1,1}+\left(t^{2}+t^{3}\right) A_{2,1,1}+A_{2,2}
\end{aligned}
$$

## k-Schur expansion of Macdonald Polynomials

(for 2 bounded partitions of 4 )

$$
\left.\left.\left.\begin{array}{rl}
H([2,2]) & q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2}
\end{array}\right]\left[\begin{array}{ccc}
0 & A_{2,1,1} & 0 \\
A_{2,2} & A_{2,1,1} & A_{1,1,1,1}
\end{array}\right]\right]\left[\begin{array}{cc}
A_{2,1,1} & 0 \\
0 & A_{1,1,1,1} \\
A_{2,2} & A_{2,1,1}
\end{array}\right]\right)
$$

(for 2 bounded partitions of 6 )

## k-Schur expansion of Macdonald Polynomials

(for 2 bounded partitions of 4 )
$\mathrm{H}([2,2]) \quad q^{2} A_{1,1,1,1}+(q+t q) A_{2,1,1}+A_{2,2} \quad\left[\begin{array}{ccc}0 & A_{2,1,1} & 0 \\ A_{2,2} & A_{2,1,1} & A_{1,1,1,1}\end{array}\right]$

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\begin{array}{r}
\mathrm{H}\left([2,1,1] \quad t q A_{1,1,1,1}+\left(q+t^{2}\right) A_{2,1,1}+A_{2,2}\left[\begin{array}{cc}
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\end{array}\right]\right. \\
\mathrm{H}([1,1,1,1]) \quad t^{4} A_{1,1,1,1}+\left(t^{2}+t^{3}\right) A_{2,1,1}+A_{2,2}
\end{array}
$$

(for 2 bounded partitions of 6 )

$$
H_{2,2,2}, \quad->, \quad,\left[\begin{array}{cccc}
0 & A_{2,2,1,1} & A_{2,1,1,1,1} & 0 \\
0 & A_{2,2,1,1} & A_{2,1,1,1,1} & 0 \\
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## Some remarkable determinants

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| $(1,0)$ | $(1,1)$ |
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If $\left(\mathrm{p}_{\mathbf{1}}, \mathrm{q}_{1}\right),\left(\mathrm{p}_{\mathbf{2}}, \mathrm{q}_{\mathbf{2}}\right), \ldots,\left(\mathrm{p}_{\mathbf{n}}, \mathrm{q}_{\mathbf{n}}\right)$ are the cells of the Ferrers diagram of $\mu \vdash \mathbf{n}$ then

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\Delta_{\mu}(x, y)=\operatorname{det}\left\|x_{j}^{p_{i}} y_{j}^{q_{i}}\right\|_{i, j=1}^{n}
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## Brief review

## Theorem (easy)

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```
\[
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```


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diff(D21,y3);
diff(D21,x3,y2);
\[
\begin{gathered}
y 3-y 2 \\
y 2-y 1 \\
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1
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$$
\begin{aligned}
& D 2 I:=\left[\begin{array}{ccc}
1 & 1 & 1 \\
y I & y 2 & y 3 \\
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\end{array}\right] \\
& y 2 x 3-y 3 x 2-y l x 3+y l x 2+x l y 3-x I y 2 \\
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$$
\begin{aligned}
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$\operatorname{diff}(\mathrm{D} 21, x 1)$;
$\operatorname{diff}(\mathrm{D} 21, x 3) ;$
$\operatorname{diff}(\mathrm{D} 21, \mathrm{y} 1) ;$
$\operatorname{diff}(\mathrm{D} 21, \mathrm{Y} 3) ;$


$$
-x 3+x 2
$$

$$
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1

## Theorem (easy)

## Brief review

For any $\mu \vdash n$ the dimension of the linear span of the derivatives of $\Delta_{\mu}(X, Y)$ is at most $n$ ! In symbols $\operatorname{dim} \mathbf{M}_{\mu}[\mathbf{X}, \mathbf{Y}] \leq \mathbf{n}!$
For example [ using MAPLE]
$\operatorname{DDmu}([2,1]) ;$
D21: $=\operatorname{det}(") ;$
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## Hilbert series

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Our spaces $\mathrm{M}[\mathrm{X}, \mathrm{Y}]$ are "bigraded" that is we have the double decomposition

$$
\mathbf{M}_{\mu}[\mathbf{X}, \mathbf{Y}]=\bigoplus_{\mathbf{r}=0}^{\mathbf{n}(\mu)} \bigoplus_{\mathrm{s}=0} \bigoplus_{\mathrm{n}}\left(\mu^{\prime}\right) \mathbf{H} \mathrm{H}_{\mathbf{r}, \mathrm{s}}\left(\mathbf{M}_{\mu}[\mathbf{X}, \mathrm{Y}]\right)
$$

With $\mathrm{H}_{\mathbf{r}, \boldsymbol{s}}\left(\mathrm{M}_{\mu}[\mathbf{X}, \mathrm{Y}]\right)$ the linear span of derivatives of $\Delta_{\mu}(\mathbf{x}, \mathbf{y})$
that are homogeneous of degree $r$ in $x_{1}, x_{2}, \ldots, x_{n}$ and degree $S$ in $y_{1}, y_{2}, \ldots, y_{n}$
Here and after we set

## Hilbert series

A vector space V is called "graded" if and only if

$$
\mathbf{V}=\mathrm{H}_{\mathbf{o}}(\mathbf{V}) \oplus \mathrm{H}_{\mathbf{1}}(\mathbf{V}) \oplus \mathrm{H}_{\mathbf{2}}(\mathbf{V}) \oplus \cdots \oplus \mathrm{H}_{\mathbf{m}}(\mathbf{V}) \oplus \cdots
$$

The subspace " $\mathrm{H}_{\mathrm{m}}(\mathrm{V})$ " is called the " m th homogeneous component" of V .
its elements are called homogeneous of degree m
If $\operatorname{dim} \mathrm{H}_{\mathrm{m}}(\mathrm{V})<\infty$ for all m , we set

$$
\mathrm{F}_{\mathrm{V}}(\mathrm{t})=\sum_{\mathrm{m}>0} \mathrm{t}^{\mathrm{m}} \operatorname{dim} \mathrm{H}_{\mathrm{m}}(\mathrm{~V})
$$

For instance for $\mathbf{R}=\mathbb{Q}\left[\mathbf{x}_{\mathbf{1}}, \mathbf{x}_{2}, \ldots, \mathbf{x}_{\mathbf{n}}\right]$ we have $\quad R=H_{0}(R) \oplus H_{1}(R) \oplus H_{2}(R) \oplus \cdots \quad$ with

$$
\mathrm{H}_{\mathrm{m}}(\mathrm{R})=\mathrm{L}\left[\mathrm{x}_{1}^{\mathrm{P} 1} \mathrm{x}_{2}^{\mathrm{P} 2} \cdots \mathrm{x}_{\mathrm{n}}^{\mathrm{Pn}}: \mathrm{p}_{1}+\mathrm{p}_{2}+\cdots+\mathrm{p}_{\mathrm{n}}=\mathrm{m}\right]
$$

In this case

$$
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$$
\mathrm{F}_{\mu}(\mathbf{q}, \mathrm{t})=\sum_{\mathrm{r}=0}^{\mathbf{n}(\mu) \mathbf{n}\left(\mu_{\mathrm{s}=\mathbf{0}}\right)} \mathrm{t}^{\mathrm{r}} \boldsymbol{q}^{\mathrm{s}} \operatorname{dim} \mathrm{H}_{\mathrm{r}, \mathrm{~s}}\left(\mathrm{M}_{\mu}[\mathbf{X}, \mathrm{Y}]\right)
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## The Macdonald Polynomials as Frobenius Characteristics

The Macdonald Polynomials as Frobenius Characteristics

$$
\mathbf{M}_{\mu}[X, Y]=\bigoplus_{r=0}^{n(\mu)} \bigoplus_{s=0}^{n\left(\mu^{\prime}\right)} \mathcal{H}_{r, s}\left(\mathbf{M}_{\mu}[X, Y]\right)
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The Frobenius map

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The Frobenius map

$$
\mathrm{F} \chi^{\lambda}=\mathbf{S}_{\lambda}[\mathbf{X}]
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The Macdonald Polynomials as Frobenius Characteristics

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It follows from the n!-Theorem that

$$
\widetilde{H}_{\mu}[X ; q, t]=\sum_{r=0}^{n(\mu)} \sum_{s=0}^{n\left(\mu^{\prime}\right)} t^{r} q^{s} \mathbf{F} \operatorname{char} \mathcal{H}_{r, s}\left(\mathbf{M}_{\mu}[X, Y]\right)
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Moreover the Frobenius characteristic of restriction of $M_{\mu}[\mathrm{X}, \mathrm{Y}]$ to $\mathrm{S}_{\mathbf{n}-\mathbf{1}}$

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\partial_{p_{1}} \widetilde{H}_{\mu}[X ; q, t]
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The Macdonald Polynomials as Frobenius Characteristics

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\widetilde{H}_{2,2}[X ; q, t]=s_{4}+s_{3,1}(q+t+q t)+s_{2,2}\left(q^{2}+t^{2}\right)+s_{2,1,1}\left(q t^{2}+q^{2} t+q t\right)+s_{1,1,1,1} t^{2} q^{2}
$$

The Macdonald Polynomials as Frobenius Characteristics

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\partial_{p_{1}} \widetilde{H}_{2,2}[X ; q, t]=(1+q+t+q t) s_{3}+\left(q^{2}+t^{2}+q t^{2}+q^{2} t+q+t+2 q t\right) s_{2,1}+q t(1+q+t+q t) s_{1,1,1}
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How does this decompose in terms of the basis $\left\{\tilde{\mathrm{H}}_{\mu}[\mathrm{X} ; \mathrm{q}, \mathrm{t}]\right\}_{\mu+3}$

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\end{array}\right] \stackrel{\text { Restriction to } \mathcal{s}_{3}}{ } \\
& \text { How does this decompose in terms of the basis }\left\{\tilde{\mathrm{H}}_{\mu}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]\right\}_{\mu+3} ? ?
\end{aligned}
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s_{3} & s_{3}+s_{2,1} & s_{2,1}
\end{array}\right] \stackrel{\text { Restriction to } \mathcal{S}_{3}}{\rightleftarrows} \\
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\end{aligned}
$$

## The Macdonald-Stanley (dual) Pieri Rules

## The Macdonald-Stanley [dual) Pieri Rules

$$
\partial_{\mathbf{p} \mathbf{1}} \tilde{\mathrm{H}}_{\mu}[\mathbf{X} ; \mathbf{q}, \mathbf{t}]=\sum_{\nu \rightarrow \mu} \mathbf{c}_{\mu, \nu}(\mathbf{q}, \mathrm{t}) \tilde{\mathrm{H}}_{\nu}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]
$$

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\begin{aligned}
& \partial_{\mathbf{p} \mathbf{1}} \tilde{H}_{\mu}[\mathbf{X} ; \mathbf{q}, \mathrm{t}]=\sum_{\nu \rightarrow \mu} \mathbf{c}_{\mu, \nu}(\mathbf{q}, \mathrm{t}) \tilde{H}_{\nu}[\mathbf{X} ; \mathbf{q}, \mathrm{t}] \\
& \mathbf{c}_{\mu \nu}(\mathbf{q}, \mathrm{t})=\prod_{\mathbf{s} \in \mathbf{R}_{\mu^{\prime}}} \frac{\mathbf{t}^{\mathbf{1}_{\mu}(\mathbf{s})}-\mathbf{q}^{\mathbf{a}_{\mu}(\mathbf{s})+\mathbf{1}}}{\mathbf{1}^{\mathbf{l}_{2}(\mathbf{s})}-\mathbf{q}^{\mathbf{a}_{\nu}(\mathbf{s})+\mathbf{1}}} \prod_{\mathbf{s} \in \mathbf{C}_{\mu^{\mu}}} \frac{\mathbf{q}^{\mathbf{a}_{\mu}(\mathbf{s})}-\mathbf{t}^{\mathbf{1}_{\mu}(\mathbf{s})+\mathbf{1}}}{\mathbf{a}^{\mathbf{a}_{2}(\mathbf{s})}-\mathbf{t}^{\mathbf{1}_{\nu}(\mathbf{s})+\mathbf{1}}}
\end{aligned}
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## The Macdonald-Stanley [dual) Pieri Rules

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F_{\mu}(q, t)=\sum_{\nu \rightarrow \mu} c_{\mu \nu}(q, t) F_{\nu}(q, t)
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hilb([3,2]);
```


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hilb([3,2]);

$$
q^{4} t^{2}+4 q^{4} t+4 q^{3} t^{2}+5 q^{4}+15 q^{3} t+9 q^{2} t^{2}+11 q^{3}+22 q^{2} t+11 q t^{2}+9 q^{2}+15 q t+5 t^{2}+4 q+4 t+1
$$

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hilb([3,2]);

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$$

$$
\left[\begin{array}{rrrrr}
5 & 11 & 9 & 4 & 1 \\
4 & 15 & 22 & 15 & 4 \\
1 & 4 & 9 & 11 & 5
\end{array}\right]
$$

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hilb([3,2]);

$$
\left.\begin{array}{rl}
q^{4} t^{2}+4 q^{4} t+4 q^{3} t^{2}+5 q^{4}+15 & q^{3} t+9 \\
2 & q^{2} t^{2}+11 \\
\hline
\end{array} q^{3}+22 q^{2} t+11 q t^{2}+9 q^{2}+15 q t+5 t^{2}+4 q+4 t+1\right)\left[\begin{array}{rrrrr}
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# The underlying representation theoretical identity 

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$\frac{\mathrm{n}!}{2}$ Conjecture
If $\alpha, \beta \vdash n$ differ only by the position of one corner cell then $\quad \operatorname{dim} \mathrm{M}_{\alpha} \wedge \mathrm{M}_{\beta}=\frac{\mathrm{n}!}{2}$
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## next the miracles

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    (3) \(\quad \mathrm{G}_{\mathrm{D}}(\mathrm{x} ; \mathbf{q}, \mathbf{t})=\mathrm{G}_{\mathrm{D}_{1}}(\mathrm{x} ; \mathbf{q}, \mathbf{t}) \mathrm{G}_{\mathrm{D}_{\mathbf{2}}}(\mathrm{x} ; \mathbf{q}, \mathrm{t}) \quad\) if \(\quad \mathrm{D} \approx \mathrm{D}_{\mathbf{1}} \times \mathrm{D}_{\mathbf{2}}\)
    (4) \(\quad \partial_{\mathrm{P}_{1}} \mathrm{G}_{\mathrm{D}}(\mathrm{x} ; \mathbf{q}, \mathbf{t})=\sum_{\mathrm{s} \in \mathrm{D}} \mathrm{w}_{\mathrm{s}, \mathrm{D}}(\mathbf{q}, \mathbf{t}) \mathrm{G}_{\mathrm{D} / \mathrm{s}}(\mathrm{x} ; \mathbf{q}, \mathbf{t})\),
```

Representation theoretical reasons suggest that,

## The basic tool : "gistol" Macdonalds

(1) A "gistol" is a lattice diagram that can be transformed to a skew diagram by row and column interchanges
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```
\(\left(\right.\) (0) \(\quad \mathrm{G}_{\mathrm{D}}(\mathbf{x} ; \mathbf{q}, \mathbf{t})=\overline{\mathbf{H}}_{\mu}(\mathrm{x} ; \mathbf{q}, \mathbf{t}) \quad\) if D is the diagram of \(\mu\)
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a) $\quad \mathrm{w}_{1}[\mathrm{~s}, \mathrm{D}]=\mathrm{t}^{\mathrm{l}_{\mathrm{D}}(\mathrm{s})} \mathrm{q}^{\mathrm{a}_{\mathrm{D}}(\mathrm{s})}$
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$$
\begin{aligned}
& \left(\begin{array}{l}
(0)
\end{array} \mathrm{G}_{\mathrm{D}}(\mathbf{x} ; \mathbf{q}, \mathbf{t})=\overline{\mathbf{H}}_{\mu}(\mathrm{x} ; \mathbf{q}, \mathbf{t}) \quad \text { if } \mathrm{D} \text { is the diagram of } \mu\right. \\
& \text { (1) } \quad \mathrm{G}_{\mathrm{D}_{1}}(\mathrm{x} ; \mathbf{q}, \mathrm{t})=\mathrm{G}_{\mathrm{D}_{2}}(\mathrm{x} ; \mathbf{q}, \mathrm{t}) \quad \text { if } \quad \mathrm{D}_{1} \approx \mathrm{D}_{2} \\
& \text { (2) } \quad \mathrm{G}_{\mathrm{D}_{1}}(\mathrm{x} ; \mathbf{q}, \mathrm{t})=\mathrm{G}_{\mathrm{D}_{2}}(\mathrm{x} ; \mathbf{t}, \mathbf{q}) \\
& \text { if } \mathrm{D}_{2} \approx \mathrm{D}_{1}^{\prime} \\
& \text { (3) } \quad \mathrm{G}_{\mathrm{D}}(\mathrm{x} ; \mathbf{q}, \mathbf{t})=\mathrm{G}_{\mathrm{D}_{1}}(\mathrm{x} ; \mathbf{q}, \mathrm{t}) \mathrm{G}_{\mathrm{D}_{2}}(\mathrm{x} ; \mathbf{q}, \mathrm{t}) \quad \text { if } \quad \mathrm{D} \approx \mathrm{D}_{\mathbf{1}} \times \mathrm{D}_{2} \\
& \text { (4) } \quad \partial_{\mathrm{P}_{1}} \mathrm{G}_{\mathrm{D}}(\mathbf{x} ; \mathbf{q}, \mathbf{t})=\sum_{\mathrm{s} \in \mathrm{D}} \mathbf{w}_{\mathrm{s}, \mathrm{D}}(\mathbf{q}, \mathbf{t}) \mathrm{G}_{\mathrm{D} / \mathrm{s}}(\mathbf{x} ; \mathbf{q}, \mathbf{t}) \text {, }
\end{aligned}
$$

Representation theoretical reasons suggest that,
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a) $\quad w_{1}[s, D]=t^{l_{D}(s)} q^{a_{D}^{\prime}(s)}$
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> existence is by no means guaranteed.

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Note: these properties overdetermine the family $\left\{\mathrm{G}_{\mathrm{D}}(\mathrm{x} ; \mathbf{q}, \mathbf{t})\right\}_{\mathrm{D}}$,

An example of a use of gistols

An example of a use of gistols

An example of a use of gistols

$\partial_{p_{1}} \square=$

An example of a use of gistols



An example of a use of gistols



An example of a use of gistols



An example of a use of gistols



An example of a use of gistols



An example of a use of gistols



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An example of a use of gistols


An example of a use of gistols

$\left.\partial_{\mathbf{p} \mathbf{1}} \square=\square+\left(\mathrm{t}+\mathrm{t}^{2}\right)^{\square}-(\mathrm{q}+\mathrm{tc})^{\square}\right)^{\square}$

An example of a use of gistols

$\partial_{\mathbf{P i}} \square=\square+\left(\mathbf{t}+\mathbf{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathbf{t q})^{\square} \square$
$\partial_{\mathbf{P i}} \square=\mathbf{t}^{2} \square+(\mathbf{q}+\mathbf{q})^{\square} \square+(1+\mathbf{t})^{\square} \square$

An example of a use of gistols

$\partial_{P i} \square_{\square}^{\square}=\square+\left(t+t^{2}\right)^{\square} \square+\left(\mathbf{Q}+\mathrm{a}+()^{\square}\right.$
$\partial_{\mathrm{P} \mathbf{1}} \square=\mathrm{t}^{2} \square+(\mathrm{q}+\mathrm{q})^{\square} \square+(1+\mathrm{t})^{\square} \square$

An example of a use of gistols



$\partial_{\mathbf{P} \mathbf{1}} \square=\square+\left(\mathbf{t}+\mathbf{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathbf{t a}) \square$


An example of a use of gistols






An example of a use of gistols

$\partial_{\mathbf{P i}} \square=\square+\left(\mathbf{t}+\mathbf{t}^{2}\right) \square+(\mathbf{q}+\mathbf{t q}) \square$
$\partial_{\mathrm{Pi}} \square=\mathrm{t}^{2} \square+(\mathbf{q}+\mathrm{q}) \square \square+(1+\mathrm{t}) \square \square=\frac{\square}{\square}=\frac{(1-\mathrm{t}) \square+(\mathrm{a}-1) \square-\square}{\mathrm{a}-\mathrm{t}}$
$\square=\square \Delta \square+\mathrm{t}^{2} \mathrm{q}^{2}+\square \Delta \square=\phi+\mathrm{q}^{\psi}$


An example of a use of gistols


An example of a use of gistols


An example of a use of gistols



$\partial_{\mathbf{P i}} \square=\square+\left(\mathbf{t}+\mathbf{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathbf{t q})^{\square} \square$

$\square=\square \Delta \square+\mathrm{t}^{2} \mathbf{q}^{2}+\square \Delta \square=\phi+\mathrm{q}^{\omega}$
$\square=\square \Delta \square+\mathrm{t}^{3} \mathrm{q} \downarrow \square \Delta \square=\phi+\mathrm{t} \dot{\square}$

$$
\square=\phi+\psi
$$



An example of a use of gistols



$\partial_{\mathbf{p} \mathbf{1}} \square=\square+\left(\mathrm{t}+\mathrm{t}^{2}\right)^{\square}-(\mathrm{q}+\mathrm{tq})^{\square}$
$\partial_{\mathbf{P} \mathbf{1}} \square=\mathbf{t}^{2} \square+(\mathbf{q}+\mathbf{q}) \square+(1+\mathrm{t}) \square \square=\frac{(1-\mathrm{t}) \square+(\mathbf{q}-1) \square}{\mathbf{q}-\mathrm{t}}$
$\exists=\boxminus \wedge \boxminus+\mathrm{t}^{2} \mathrm{c}^{2} \downarrow \boxminus \wedge \boxminus=\phi+\mathrm{q} \psi$
$\boxminus=\boxminus \wedge \theta+\mathrm{t}^{3} \mathrm{q}+\boxminus \wedge \boxminus=\phi+\mathrm{t} \psi$ $\square=\phi+\psi$

$\partial_{\mathbf{P} \mathbf{i}} \boxminus=\mathbf{t}^{2}(\phi+\mathbf{q} \psi)+(\mathbf{q}+\mathbf{q} \mathbf{t})(\phi+\psi)+(1+\mathrm{t})(\phi+\mathrm{t} \psi)$

An example of a use of gistols

$\partial_{\mathbf{P i}} \square=\square+\left(\mathbf{t}+\mathbf{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathbf{t c})^{\square} \square$

$\square=\square \Delta \square+\mathrm{t}^{2} \mathbf{q}^{2}+\square \Delta \square=\phi+\mathbf{q}^{\omega}$
$\square=\square \Delta \square+\mathrm{t}^{3} \mathrm{a}+\square \Delta \square=\phi+\mathrm{t} \dot{\square}$

$\partial_{\mathbf{P} \mathbf{1}} \square^{\square}=\mathbf{t}^{2}(\phi+\mathbf{q} \psi)+(\mathbf{q}+\mathbf{q})(\phi+\psi)+(1+\mathbf{t})(\phi+\mathbf{t} \psi)$

$$
=\left(t^{2}+\mathbf{q}+\mathbf{q}+1+\mathrm{t}\right) \phi+\left(\mathrm{t}^{2} \mathrm{a}+\mathrm{q}+\mathrm{q} \mathrm{t}+\mathrm{t}+\mathrm{t}^{2}\right) \phi
$$

An example of a use of gistols

$\partial_{\mathbf{P i}} \square=\square+\left(\mathrm{t}+\mathrm{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathbf{t a})^{\square} \square$

$\square=\square \Delta \square+\mathrm{t}^{2} \mathbf{q}^{2}+\square \Delta \square=\phi+\mathbf{q}^{\omega}$
$\square=\square \Delta \square+\mathrm{t}^{3} \mathrm{a}+\square \Delta \square=\phi+\mathrm{t} \dot{\square}$


$$
\begin{array}{r}
\partial_{\mathbf{p} \mathbf{t}} \boxminus=\mathrm{t}^{2}(\phi+\mathbf{q} \psi)+(\mathbf{q}+\mathbf{q t})(\phi+\psi)+(1+\mathrm{t})(\phi+\mathrm{t} \psi) \\
=\left(\mathrm{t}^{2}+\mathbf{q}+\mathbf{q} \mathbf{t}+1+\mathrm{t}\right) \phi+\left(\mathrm{t}^{2} \mathbf{q}+\mathbf{q}+\mathbf{q} \mathbf{t}+\mathrm{t}+\mathrm{t}^{2}\right) \psi \\
=\left((1+\mathrm{t})(1+\mathbf{q})+\mathrm{t}^{2}\right) \phi+(\mathrm{t}(1+\mathrm{t})(1+\mathbf{q})+\mathbf{q}) \psi
\end{array}
$$

An example of a use of gistols

$\partial_{\mathbf{P i}} \square=\square+\left(\mathrm{t}+\mathrm{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathbf{t a})^{\square} \square$

$\square=\square \Delta \square+\mathrm{t}^{2} \mathbf{q}^{2}+\square \Delta \square=\phi+\mathbf{q}^{\omega}$
$\square=\square \Delta \square+\mathrm{t}^{3} \mathrm{a}+\square \Delta \square=\phi+\mathrm{t} \dot{\square}$


$$
\begin{array}{r}
\partial_{\mathbf{p} \mathbf{t}} \boxminus=\mathrm{t}^{2}(\phi+\mathbf{q} \psi)+(\mathbf{q}+\mathbf{q t})(\phi+\psi)+(1+\mathrm{t})(\phi+\mathrm{t} \psi) \\
=\left(\mathrm{t}^{2}+\mathbf{q}+\mathbf{q} \mathbf{t}+1+\mathrm{t}\right) \phi+\left(\mathrm{t}^{2} \mathbf{q}+\mathbf{q}+\mathbf{q} \mathbf{t}+\mathrm{t}+\mathrm{t}^{2}\right) \psi \\
=\left((1+\mathrm{t})(1+\mathbf{q})+\mathrm{t}^{2}\right) \phi+(\mathrm{t}(1+\mathrm{t})(1+\mathbf{q})+\mathbf{q}) \psi \\
=(1+\mathrm{t})(1+\mathbf{q}) \operatorname{t}+\mathrm{t}^{2} \phi+\mathbf{q} \psi
\end{array}
$$

An example of a use of gistols

\section*{|  |  |  |  |
| :--- | :--- | :--- | :--- |
|  |  |  |  |
| - |  | $A$ | $t^{l}$ |$q^{a^{\prime}}$}




$\partial_{\mathbf{p} 1} \square=\square+\left(\mathrm{t}+\mathrm{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathrm{tq})^{\square}$
$\partial_{\mathrm{P} \mathbf{1}} \square=\mathrm{t}^{2} \square+(\mathrm{q}+\mathrm{qt}) \square+(1+\mathrm{t}) \square \mathrm{\square} \square=\frac{(1-\mathrm{t}) \square+(\mathrm{q}-1) \square}{\mathrm{q}-\mathrm{t}}$
$\boxminus=\boxminus \wedge \boxminus+\mathrm{t}^{2} \mathrm{c}^{2} \downarrow \boxminus \wedge \boxminus=\phi+\mathrm{q} \psi$ $\boxminus=\boxminus \wedge \boxminus+\mathrm{t}^{3} \mathrm{q}+\boxminus \wedge \boxminus=\phi+\mathrm{t} \psi$

$\partial_{\mathbf{P} \mathbf{1}} \square=\mathrm{t}^{2}(\phi+\mathrm{q} \psi)+(\mathbf{q}+\mathbf{q} \mathrm{t})(\phi+\psi)+(1+\mathrm{t})(\phi+\mathrm{t} \psi)$

$$
\begin{gathered}
=\left(\mathrm{t}^{2}+\mathrm{q}+\mathrm{q} \mathrm{t}+1+\mathrm{t}\right) \phi+\left(\mathrm{t}^{2} \mathrm{q}+\mathrm{q}+\mathrm{q} \mathrm{t}+\mathrm{t}+\mathrm{t}^{2}\right) \psi \\
=\left((1+\mathrm{t})(1+\mathrm{q})+\mathrm{t}^{2}\right) \phi+(\mathrm{t}(1+\mathrm{t})(1+\mathrm{q})+\mathrm{q}) \psi \\
=(1+\mathrm{t})(1+\mathrm{q}) \\
=+\mathrm{t}^{2} \phi+\mathrm{q} \psi
\end{gathered}
$$

An example of a use of gistols


$\partial_{\mathbf{P i}}{ }^{\square}=\boxminus+\left(\mathbf{t}+\mathrm{t}^{2}\right)^{\square} \square+(\mathbf{q}+\mathrm{tq})^{\square}$

$\exists=\boxminus \wedge \boxminus+\mathrm{t}^{2} \mathrm{c}^{2} \downarrow \boxminus \wedge \boxminus=\phi+\mathrm{q} \psi$
$\boxminus=\boxminus \wedge \boxminus+\mathrm{t}^{3} \mathrm{q}+\boxminus \wedge \boxminus=\phi+\mathrm{t} \psi$


$$
\begin{array}{r}
\partial_{\mathbf{P} \mathbf{1}} \square=\mathrm{t}^{2}(\phi+\mathbf{q} \psi)+(\mathbf{q}+\mathbf{q} \mathrm{t})(\phi+\psi)+(1+\mathrm{t})(\phi+\mathrm{t} \psi) \\
=\left(\mathrm{t}^{2}+\mathbf{q}+\mathbf{q}+1+\mathbf{t}\right) \phi+\left(\mathrm{t}^{2} \mathbf{q}+\mathbf{q}+\mathbf{q} \mathbf{t}+\mathrm{t}+\mathrm{t}^{2}\right) \psi \\
=\left((1+\mathrm{t})(1+\mathbf{q})+\mathrm{t}^{2}\right) \phi+(\mathrm{t}(1+\mathrm{t})(1+\mathbf{q})+\mathbf{q}) \psi \\
=(1+\mathrm{t})(1+\mathbf{q})-\left(\mathrm{t}^{2} \phi+\mathbf{q}^{\psi} \psi\right.
\end{array}
$$



An example of a use of gistols

\section*{|  |  |  |  |
| :--- | :--- | :--- | :--- |
|  |  |  |  |
| - |  | $A$ | $t^{l}$ |$q^{a^{\prime}}$}

 $+a t$



$\square=\square=\square+\square+t^{2} \mathbf{c}^{2}+\square \square=\square \square=\phi+\square$


$\partial_{\mathbf{P} \mathbf{1}} \boxminus=\mathrm{t}^{2}(\phi+\mathbf{q} \psi)+(\mathbf{q}+\mathbf{q} \mathrm{t})(\phi+\psi)+(1+\mathrm{t})(\phi+\mathrm{t} \psi)$

$$
\begin{aligned}
& =\left(\mathrm{t}^{2}+\mathrm{q}+\mathrm{q} \mathrm{t}+1+\mathrm{t}\right) \phi+\left(\mathrm{t}^{2} \mathrm{q}+\mathrm{q}+\mathrm{qt}+\mathrm{t}+\mathrm{t}^{2}\right) \psi \\
& =\left((1+\mathrm{t})(1+\mathrm{q})+\mathrm{t}^{2}\right) \phi+(\mathrm{t}(1+\mathrm{t})(1+\mathrm{q})+\mathrm{q}) \psi
\end{aligned}
$$

An example of a use of gistols

$\partial_{\mathbf{P i}} \square=\square+\left(\mathbf{t}+\mathbf{t}^{2}\right)^{\square} \square+(\mathbf{a}+\mathbf{t c})^{\square} \square$
$\partial_{\mathbf{P i}} \square=\mathrm{t}^{2} \square+(\mathbf{q}+\mathrm{q})^{\square} \square+(1+\mathrm{t}) \square \square=\frac{(1-\mathrm{t}) \square+(\mathrm{a}-1) \square-\square}{\mathrm{a}}=\frac{\square-\mathrm{a}}{\square}$
$\square=\square \Delta \square+\mathrm{t}^{2} \mathbf{q}^{2}+\square \Delta \square=\phi+\mathbf{q}^{\omega}$ $\square=\square \Delta \square+\mathrm{t}^{3} \mathrm{a}+\square \Delta \square=\phi+\mathrm{t} \dot{\square}$


$$
\begin{aligned}
& \partial_{\mathbf{P i}} \square=\mathbf{t}^{2}(\phi+\mathbf{q} \psi)+(\mathbf{q}+\mathbf{q} \mathbf{t})(\phi+\psi)+(1+\mathbf{t})(\phi+\mathbf{t} \psi) \\
& =\left(\mathrm{t}^{2}+\mathbf{q}+\mathbf{q} t+1+\mathrm{t}\right) \phi+\left(\mathrm{t}^{2} \mathrm{a}+\mathrm{q}+\mathrm{q} \mathrm{t}+\mathrm{t}+\mathrm{t}^{2}\right) \phi \\
& =\left((1+\mathrm{t})(1+\mathrm{q})+\mathrm{t}^{2}\right) \phi+(\mathrm{t}(1+\mathrm{t})(1+\mathrm{q})+\mathrm{q}) \psi
\end{aligned}
$$



An example of a use of gistols

$\partial_{\mathbf{P i}} \square=\square+\left(\mathbf{t}+\mathbf{t}^{2}\right)^{\square} \square+(\mathbf{a}+\mathbf{t c})^{\square} \square$

$\square=\square \Delta \square+\mathrm{t}^{2} \mathbf{q}^{2}+\square \Delta \square=\phi+\mathbf{q}^{\omega}$ $\square=\square \Delta \square+\mathrm{t}^{3} \mathrm{a}+\square \Delta \square=\phi+\mathrm{t} \dot{\square}$

$\partial_{\mathbf{P i}} \square=\mathbf{t}^{2}(\phi+\mathbf{c} \psi)+(\mathbf{c}+\mathbf{q} \mathbf{t})(\phi+\psi)+(1+\mathbf{t})(\phi+\mathbf{t} \psi)$

$$
\begin{aligned}
& =\left(\mathrm{t}^{2}+\mathrm{q}+\mathrm{q} \mathrm{t}+1+\mathrm{t}\right) \phi+\left(\mathrm{t}^{2} \mathrm{q}+\mathrm{q}+\mathrm{q} \mathrm{t}+\mathrm{t}+\mathrm{t}^{2}\right) \psi \\
& =\left((1+\mathrm{t})(1+\mathrm{q})+\mathrm{t}^{2}\right) \phi+(\mathrm{t}(1+\mathrm{t})(1+\mathrm{q})+\mathrm{q}) \psi
\end{aligned}
$$

$$
\square=(1+t)(1+q) \square+t^{2} \phi+q \omega=(1+t)(1+q) \square+t^{2} \square(\mathrm{q} / \mathrm{t}, \mathrm{t})
$$

## The General case

## The General case



## The General case



## The General case




## The General case



## The General case



## The General case




## The General case



## The General case



## The General case



## The General case

$\because=\phi_{\mathrm{ab}}+\mathrm{t}^{\mathrm{a}-\mathrm{b}} \psi_{\mathrm{ab}}$


## The General case



## The General case



## The General case



## The General case



## The General case



## The General case



## The General case



## The General case



$$
\begin{aligned}
& \stackrel{\square}{\square}{ }_{\partial_{\mathbf{P}}}^{\square}=\mathbf{t}^{\mathbf{b}}[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)+\mathbf{q}[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a - b}} \psi_{\mathbf{a b}}\right) \\
& \psi_{\mathrm{ab}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t}
\end{aligned}
$$

## The General case



$$
\begin{aligned}
& \left.\partial_{\mathbf{p} 1}{ }^{\square}=\mathrm{t}^{\mathbf{b}}[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)+\mathrm{q}^{[\mathrm{b}}\right]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a - b}} \psi_{\mathbf{a b}}\right) \\
& \begin{array}{c}
\psi_{\mathrm{ab}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}}
\end{array}
\end{aligned}
$$

## The General case




$$
\begin{array}{r}
\partial_{\mathbf{p} \mathbf{t}}=\mathrm{t}^{\mathbf{b}}[\mathbf{a}-\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)+\mathbf{q}^{[\mathrm{b}} \mathbf{t}_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a - b}} \psi_{\mathbf{a b}}\right) \\
\psi_{\mathbf{a b}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}} \\
\psi_{\mathbf{a b}}: q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t}
\end{array}
$$

## The General case




$$
\begin{array}{r}
\partial_{\mathbf{p} \mathbf{t}}=\mathrm{t}^{\mathbf{b}}[\mathbf{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)+\mathbf{q}^{[\mathrm{b}} \mathbf{t}_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a - b}} \psi_{\mathbf{a b}}\right) \\
\psi_{\mathbf{a b}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}} \\
\psi_{\mathbf{a b}}: q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t} \\
\psi_{\mathbf{a b}}: q[a-b]_{t}+(1+q) t^{a-b}[b]_{t}
\end{array}
$$

## The General case



$\partial_{\mathbf{P} \mathbf{t}}^{\square}=\mathrm{t}^{\mathrm{b}}[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)+\mathbf{q}[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a - b}} \psi_{\mathbf{a b}}\right)$

$$
\phi_{\mathrm{ab}}: t^{b}[a-b]_{t}+(1+q)[b]_{t}
$$

$$
\begin{aligned}
\psi_{\mathrm{ab}} & : q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t} } & =1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t} \\
\psi_{\mathrm{ab}} & : q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t} \\
\psi_{\mathbf{a b}} & : q[a-b]_{t}+(1+q) t^{a-b}[b]_{t}
\end{aligned}
$$

## The General case



$=\phi_{\mathrm{ab}}+\psi_{\mathrm{ab}}$
$\begin{array}{r}\square \\ \partial_{\mathbf{p} \mathbf{1}} \\ \square\end{array}=\mathrm{t}^{\mathbf{b}}[\mathbf{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)+\mathbf{q}[\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)$

$$
\phi_{\mathrm{ab}}: t^{b}[a-b]_{t}+(1+q)[b]_{t}
$$

$$
\begin{aligned}
\psi_{\mathrm{ab}} & : q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t} } & =1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t} \\
\psi_{\mathrm{ab}} & : q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t} \\
\psi_{\mathrm{ab}} & : q[a-b]_{t}+(1+q) t^{a-b}[b]_{t}
\end{aligned}
$$

$$
\partial_{\mathbf{p}_{1}} \stackrel{B}{B}_{\square}^{B}=(1+\mathbf{q})[\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)+[\mathbf{a}-\mathbf{b}]_{\mathbf{t}}\left(\mathbf{t}^{\mathbf{b}} \phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)
$$

## The General case



$=\phi_{\mathrm{ab}}+\psi_{\mathrm{ab}}$
$\begin{array}{r}\square \\ \partial_{\mathbf{p} \mathbf{1}} \\ \square\end{array}=\mathrm{t}^{\mathbf{b}}[\mathbf{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)+\mathbf{q}[\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)$

$$
\begin{gathered}
\psi_{\mathrm{ab}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}} \\
\psi_{\mathrm{ab}}: q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t}
\end{gathered}
$$

$$
\phi_{\mathrm{ab}}: t^{b}[a-b]_{t}+(1+q)[b]_{t} \quad \quad \psi_{\mathrm{ab}}: q[a-b]_{t}+(1+q) t^{a-b}[b]_{t}
$$

$$
\partial_{\mathbf{p} 1} \boxminus=(1+\mathrm{q})[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathrm{ab}}+\mathrm{t}^{\mathbf{a - b}} \psi_{\mathbf{a b}}\right)+[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\mathrm{t}^{\mathbf{b}} \phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)
$$

## The General case



$=\mathrm{t}^{\mathrm{b}}[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{q} \psi_{\mathbf{a b}}\right)+\mathbf{q}[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathbf{a b}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)$

$$
\begin{gathered}
\psi_{\mathrm{ab}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}} \\
\psi_{\mathrm{ab}}: q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t}
\end{gathered}
$$

$$
\phi_{\mathrm{ab}}: t^{b}[a-b]_{t}+(1+q)[b]_{t} \quad \quad \psi_{\mathrm{ab}}: q[a-b]_{t}+(1+q) t^{a-b}[b]_{t}
$$



## The General case



$=\mathrm{t}^{\mathrm{b}}[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathrm{ab}}+\mathrm{q} \psi_{\mathbf{a b}}\right)+\mathrm{q}[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathrm{ab}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)$

$$
\begin{array}{r}
\psi_{\mathbf{a b}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
\phi_{\mathbf{a b}}: t^{t^{b}[a-b]_{t}+(1+q)[b]_{t}} \begin{array}{r}
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}} \\
\psi_{\mathbf{a b}}: q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t} \\
\psi_{\mathbf{a b}}: q[a-b]_{t}+(1+q) t^{a-b}[b]_{t}
\end{array} \\
\partial_{\partial_{\mathbf{p} 1}}=(1+\mathbf{q})[\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{t}^{\mathbf{a - b}} \psi_{\mathbf{a b}}\right)+[\mathbf{a - b}]_{\mathbf{t}}\left(\mathbf{t}^{\mathbf{b}} \phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)
\end{array}
$$

## The General case



$=\mathrm{t}^{\mathrm{b}}[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathrm{ab}}+\mathrm{q} \psi_{\mathbf{a b}}\right)+\mathrm{q}[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathrm{ab}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)$

$$
\begin{gathered}
\psi_{\mathrm{ab}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}} \\
\psi_{\mathrm{ab}}: q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t}
\end{gathered}
$$

$$
\phi_{\mathrm{ab}}: t^{b}[a-b]_{t}+(1+q)[b]_{t} \quad \psi_{\mathrm{ab}}: q[a-b]_{t}+(1+q) t^{a-b}[b]_{t}
$$

$$
\partial_{\mathbf{p} \mathbf{1}}{ }_{B}^{B}=(1+\mathrm{q})[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)+[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\mathrm{t}^{\mathbf{b}} \phi_{\mathbf{a b}}+\mathrm{q} \psi_{\mathbf{a b}}\right)
$$

## The General case



$=\mathrm{t}^{\mathrm{b}}[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathrm{ab}}+\mathrm{q} \psi_{\mathbf{a b}}\right)+\mathrm{q}[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\psi_{\mathrm{ab}}\right)+[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathrm{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)$

$$
\begin{gathered}
\psi_{\mathrm{ab}}: q\left(t^{b}[a-b]_{t}+[b]_{t}\right)+t^{a-b}[b]_{t} \\
{[b]_{t}+t^{b}[a-b]_{t}=1+t+\cdots+t^{a-1}=[a-b]_{t}+t^{a-b}[b]_{t}} \\
\psi_{\mathrm{ab}}: q\left([a-b]_{t}+t^{a-b}[b]_{t}\right)+t^{a-b}[b]_{t}
\end{gathered}
$$

$$
\left.\phi_{\mathrm{ab}}: t^{b}[a-b]_{t}+(1+q)[b]_{t} \quad \psi_{\mathrm{ab}}: q[a-b]_{t}+1+q\right) t^{a-b}[b]_{t}
$$

$$
\partial_{\mathbf{p}_{1}} \exists^{B}=(1+\mathbf{q})[\mathrm{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)+[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\mathbf{t}^{\mathbf{b}} \phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)
$$

## the next step

## the next step



## the next step



While my representation theoretical extension of the Haglund conjecture is

## the next step

$$
{ }_{\partial_{\mathbf{p}}} B=(1+\mathbf{q})[\mathbf{b}]_{\mathbf{t}}\left(\phi_{\mathbf{a b}}+\mathbf{t}^{\mathbf{a}-\mathbf{b}} \psi_{\mathbf{a b}}\right)+[\mathrm{a}-\mathrm{b}]_{\mathbf{t}}\left(\mathrm{t}^{\mathbf{b}} \phi_{\mathbf{a b}}+\mathbf{q} \psi_{\mathbf{a b}}\right)
$$

While my representation theoretical extension of the Haglund conjecture is


## the next step



While my representation theoretical extension of the Haglund conjecture is


## the next step



While my representation theoretical extension of the Haglund conjecture is


## the next step



While my representation theoretical extension of the Haglund conjecture is

so we are reduced to proving the identity

## the next step



While my representation theoretical extension of the Haglund conjecture is

so we are reduced to proving the identity


## the next step



While my representation theoretical extension of the Haglund conjecture is

so we are reduced to proving the identity

$k$-schur visualization of the $a=3 b=2$ case
$k$-schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathrm{p} 1} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \exists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

## $k$-schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathrm{p} \mathbf{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \exists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

## $k$-schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} \mathbf{1}} \boxplus_{(\mathbf{q}, \mathrm{t})}=(1+\mathbf{t})(1+\mathbf{q}) \boxplus_{(\mathbf{q}, \mathbf{t})}+\mathbf{t}^{2} \boxplus_{(\mathbf{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\text { 母 }=\phi_{32}+\mathrm{t} \psi_{32}
$$

## $k$-schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} \mathbf{1}} \boxplus_{(\mathbf{q}, \mathbf{t})}=(1+\mathbf{t})(1+\mathbf{q}) \boxplus_{(\mathbf{q}, \mathbf{t})}+\mathbf{t}^{2} \boxplus_{(\mathbf{q} / \mathrm{t}, \mathbf{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

母 $=\phi_{32}+\mathrm{t} \psi_{32}$
$\boxplus=\phi_{32}+\mathbf{q} \psi_{s 2}$

## $k$-schur visualization of the $a=3 b=2$ case



$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\boxminus=\phi_{32}+\mathrm{t} \psi_{32}
$$

$$
\theta=\phi_{32}+\mathrm{q} \psi_{32}
$$

Now in terms of 2 -Schur we get

## $k$-schur visualization of the $a=3 b=2$ case



$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\boxminus=\phi_{32}+\mathrm{t} \psi_{32}
$$

$$
\boxminus=\phi_{32}+\mathrm{q} \psi_{32}
$$

Now in terms of 2 -Schur we get

$$
\phi_{32}=\boxminus \wedge \boxminus=\frac{\mathrm{q} \boxminus-\mathrm{t} \boxminus}{\mathrm{q}-\mathrm{t}}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} 1} \boxplus_{(\mathbf{q}, \mathbf{t})}=(1+\mathrm{t})(1+\mathrm{q}) \exists_{(\mathbf{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{\mathrm{s} 2}+\mathbf{q} \psi_{\mathrm{s} 2}
$$

$$
母=\phi_{32}+\mathbf{t} \psi_{32} \quad \boxplus=\phi_{32}+\mathbf{q} \psi_{32}
$$

Now in terms of 2 －Schur we get

$$
\phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{22}+\mathrm{q}^{\mathrm{A}} \mathbf{A}_{21}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} 1} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) ظ_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathbf{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{\mathrm{s} 2}+\mathbf{q} \psi_{\mathrm{s} 2}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathbf{q} \psi_{s 2}
$$

Now in terms of 2 －Schur we get

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxminus=\frac{\mathrm{q} 母-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{22}+\mathrm{q}^{\mathrm{A}} \mathbf{A}_{211} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{211} & 0
\end{array}\right]
\end{aligned}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} 1} \boxplus_{(\mathbf{q}, \mathbf{t})}=(1+\mathrm{t})(1+\mathbf{q}) \boxplus_{(\mathbf{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathbf{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathbf{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{\mathrm{s} 2}+\mathbf{q} \psi_{\mathrm{s} 2}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathbf{q} \psi_{s 2}
$$

Now in terms of 2 －Schur we get

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q} \hbar-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q A}_{\mathbf{2 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{aligned}
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} 1} \boxplus_{(\mathbf{q}, \mathbf{t})}=(1+\mathrm{t})(1+\mathbf{q}) \exists_{(\mathbf{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{\mathrm{s} 2}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

Now in terms of 2 －Schur we get

$$
\begin{aligned}
\phi_{32} & =母 \wedge \boxplus=\frac{\mathrm{q} 母-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{22}+\mathrm{q}_{\mathbf{2 1 1}} \\
\phi_{32} & =\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{211} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{aligned}
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathrm{p} \mathbf{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \exists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathrm{q} \psi_{s 2}
$$

Now in terms of 2 －Schur we get

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q}_{\mathbf{A} \mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathrm{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right]
\end{aligned}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathrm{p} \mathbf{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \exists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathrm{q} \psi_{32}
$$

Now in terms of 2 －Schur we get

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q A}_{\mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{1111} & 0
\end{array}\right]
\end{aligned}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathrm{p} \mathbf{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \exists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

$$
\text { 母 }=\phi_{32}+\mathrm{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathbf{q} \psi_{32}
$$

Now in terms of 2 －Schur we get

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q}_{\mathbf{A}} \mathbf{A}_{\mathbf{2 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathrm{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{1111} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} \mathbf{1}} \boxplus_{(\mathbf{q}, \mathbf{t})}=(1+\mathbf{t})(1+\mathbf{q}) \boxplus_{(\mathbf{q}, \mathbf{t})}+\mathbf{t}^{2} \boxplus_{(\mathbf{q} / \mathrm{t}, \mathbf{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+q \psi_{32}
$$

Now in terms of 2 －Schur we get
thus

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q}_{\mathbf{A} \mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathrm{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{1111} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p i}_{1}} \nexists_{(\mathrm{q}, \mathrm{t})}=(1+\mathbf{t})(1+\mathrm{q}) \nexists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

$$
\mathrm{B}_{1}=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathbf{q} \psi_{32}
$$

Now in terms of 2 －Schur we get
thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{qt} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{\mathbf{2 1 1}} & \mathbf{A}_{\mathbf{1 1 1 1}}
\end{array}\right]
$$

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q}_{\mathbf{A}}^{\mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{211} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{1111} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k$－schur visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p},} \boxplus_{(\mathbf{q}, \mathrm{t})}=(1+\mathbf{t})(1+\mathrm{q}) \boxplus_{(\mathbf{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{\mathrm{s} 2}
$$

$$
母^{B}=\phi_{32}+\mathrm{t} \psi_{32} \quad \boxplus=\phi_{32}+\mathrm{q} \psi_{32}
$$

Now in terms of 2 －Schur we get

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q}_{\mathbf{A}}^{\mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{1111} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

thus
and thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q}^{\mathbf{A}_{211}}+\mathrm{qt} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{\mathbf{2 1 1}} & \mathbf{A}_{\mathbf{1 1 1 1}}
\end{array}\right]
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p l}_{1}} \nexists_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \nexists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\xi=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+q \psi_{32}
$$

Now in terms of 2 -Schur we get
thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q}^{\mathbf{A}_{211}}+\mathrm{qt} \mathbf{A}_{211}+\mathrm{q}^{2} \mathbf{A}_{1111}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{211} & \mathbf{A}_{1111}
\end{array}\right]
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \mathbf{A}_{22}+\mathrm{tq} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{t}^{2} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
\mathrm{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
$$

$$
\begin{aligned}
& \phi_{32}=母_{\wedge} \oplus=\frac{\mathrm{q}^{\boxminus}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q} \mathbf{A}_{\mathbf{2 1 1}} \quad \psi_{32}=\frac{\boxplus-\boxminus}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{\mathbf{1 1 1 1}} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p} \mathbf{1}} \boxplus_{(\mathbf{q}, \mathbf{t})}=(1+\mathbf{t})(1+\mathbf{q}) \boxplus_{(\mathbf{q}, \mathbf{t})}+\mathbf{t}^{2} \boxplus_{(\mathbf{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+q \psi_{32}
$$

Now in terms of 2 －Schur we get
thus
and thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q}^{\mathbf{A}} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{qt} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{\mathbf{2 1 1}} & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & \mathbf{A}_{\mathbf{1 1 1 1}}
\end{array}\right]
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \mathbf{A}_{22}+\mathrm{tq} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{t}^{2} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
\mathrm{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
$$

$$
\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\begin{aligned}
& \phi_{32}=\boxminus \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q} \mathbf{A}_{\mathbf{2 1 1}} \quad \psi_{32}=\frac{\boxplus-母}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{\mathbf{1 1 1 1}} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p r}_{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \nexists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\xi=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathbf{q} \psi_{32}
$$

Now in terms of 2 -Schur we get
thus
and thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q}^{\mathbf{A}} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{qt} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{\mathbf{2 1 1}} & \mathbf{A}_{\mathbf{1 1 1 1}}
\end{array}\right]
$$

$$
\mathrm{t}^{2} \boxminus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \mathbf{A}_{22}+\mathrm{tq} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{t}^{2} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{1111}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
$$

$$
\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}=\mathrm{t}^{2}\left(\mathbf{A}_{22}+\mathrm{q}_{\mathbf{2 1 1}}\right)+\mathrm{q}\left(\mathrm{t} \mathbf{A}_{211}+\mathrm{q}^{1} \mathbf{A}_{1111}\right)
$$

$$
\begin{aligned}
& \phi_{32}=母_{\wedge} \oplus=\frac{\mathrm{q}^{\boxminus}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q} \mathbf{A}_{\mathbf{2 1 1}} \quad \psi_{32}=\frac{\boxplus-\boxminus}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{1111} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p r}_{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \nexists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathbf{q} \psi_{32}
$$

Now in terms of 2 －Schur we get
thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q}^{\mathbf{A}_{211}}+\mathrm{q} \mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{211} & \mathbf{A}_{1111}
\end{array}\right]
$$

$$
\begin{aligned}
& \text { and thus } \\
& \qquad \mathrm{t}^{2} \boxminus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \mathbf{A}_{22}+\mathrm{tq} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{t}^{2} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right] \\
& \mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}=\mathrm{t}^{2}\left(\mathbf{A}_{\mathbf{2 2}}+\mathrm{q}_{\mathbf{2 1 1}}\right)+\mathrm{q}\left(\mathrm{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q} \mathbf{A}_{\mathbf{1 1 1 1}}\right)=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{\mathbf{1 1 1 1}}
\end{array}\right]
\end{aligned}
$$

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q A}_{\mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{\mathbf{1 1 1 1}} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p r}_{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \nexists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+\mathbf{q} \psi_{32}
$$

Now in terms of 2 －Schur we get
thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q}^{\mathbf{A}_{211}}+\mathrm{q} \mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{211} & \mathbf{A}_{1111}
\end{array}\right]
$$

$$
\begin{aligned}
& \text { and thus } \\
& \qquad \mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \mathbf{A}_{22}+\mathrm{tq} \mathbf{A}_{211}+\mathrm{t}^{2} \mathbf{A}_{211}+\mathrm{q}^{2} \mathbf{A}_{1111}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right] \\
& \mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}=\mathrm{t}^{2}\left(\mathbf{A}_{22}+\mathrm{q}_{\mathbf{2 1 1}}\right)+\mathrm{q}\left(\mathrm{t} \mathbf{A}_{211}+\mathrm{q} \mathbf{A}_{1111}\right)=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q A}_{\mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
0 & \mathbf{A}_{\mathbf{1 1 1 1}} & 0
\end{array}\right] \quad \mathrm{q} \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right]
\end{aligned}
$$

## $k-s c h u r$ visualization of the $a=3 b=2$ case

$$
\partial_{\mathbf{p r}_{1}} \boxplus_{(\mathrm{q}, \mathrm{t})}=(1+\mathrm{t})(1+\mathrm{q}) \nexists_{(\mathrm{q}, \mathrm{t})}+\mathrm{t}^{2} \boxplus_{(\mathrm{q} / \mathrm{t}, \mathrm{t})}
$$

$$
\mathrm{t}^{2} \boxplus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}
$$

$$
\text { 母 }=\phi_{32}+\mathbf{t} \psi_{32}
$$

$$
\boxplus=\phi_{32}+q \psi_{32}
$$

Now in terms of 2 －Schur we get
thus

$$
\boxminus=\mathbf{A}_{22}+\mathrm{q}^{\mathbf{A}_{211}}+\mathrm{q} \mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & \mathbf{A}_{211} & 0 \\
\mathbf{A}_{22} & \mathbf{A}_{211} & \mathbf{A}_{1111}
\end{array}\right]
$$

$$
\begin{aligned}
& \text { and thus } \\
& \qquad \mathrm{t}^{2} \boxminus(\mathrm{q} / \mathrm{t}, \mathrm{t})=\mathrm{t}^{2} \mathbf{A}_{22}+\mathrm{tq} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{t}^{2} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q}^{2} \mathbf{A}_{\mathbf{1 1 1 1}}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & \mathbf{A}_{211} & 0 \\
0 & 0 & \mathbf{A}_{1111}
\end{array}\right] \\
& \mathrm{t}^{2} \phi_{32}+\mathrm{q} \psi_{32}=\mathrm{t}^{2}\left(\mathbf{A}_{22}+\mathrm{q}_{\mathbf{2 1 1}}\right)+\mathrm{q}\left(\mathrm{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathrm{q} \mathbf{A}_{\mathbf{1 1 1 1}}\right)=\left[\begin{array}{ccc}
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\end{aligned}
$$

$$
\begin{aligned}
& \phi_{32}=母 \wedge \boxplus=\frac{\mathrm{q}^{母}-\mathrm{t} \boxplus}{\mathrm{q}-\mathrm{t}}=\mathbf{A}_{\mathbf{2 2}}+\mathbf{q A}_{\mathbf{2 1 1}} \quad \psi_{32}=\frac{母-母}{\mathrm{q}-\mathrm{t}}=\mathbf{t} \mathbf{A}_{\mathbf{2 1 1}}+\mathbf{q} \mathbf{A}_{\mathbf{1 1 1 1}} \\
& \phi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
0 & 0 & 0 \\
\mathbf{A}_{\mathbf{2 2}} & \mathbf{A}_{\mathbf{2 1 1}} & 0
\end{array}\right] \quad \mathrm{t}^{2} \phi_{32}=\left[\begin{array}{ccc}
\mathbf{A}_{22} & \mathbf{A}_{211} & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \quad \psi_{32}=\left[\begin{array}{ccc}
0 & 0 & 0 \\
\mathbf{A}_{211} & 0 & 0 \\
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## A simpler question?

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Count the number of permutations of $S_{n}$

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\sum_{a \in C_{n, k}} \operatorname{sign}(a)=\sum_{r=0}^{k}\binom{k}{r}(-1)^{k-r} n(n-1) \cdots(n-r+1)
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$$



Do you think that was enough?

Do you think that was enough? no!

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here is more

## Do you think that was enough? no! here is more

$$
\sum_{a \in \Pi_{n, k}} q^{\operatorname{maj}\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[n]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q}
$$

## Do you think that was enough? no! here is more

$$
\sum_{a \in \Pi_{n, k}} q^{m a j\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[n]_{\mathbf{q}}[n-1]_{q} \cdots[n-r+1] \mathbf{q}
$$

Proof (sketch)

## Do you think that was enough? no! here is more

$$
\sum_{\mathbf{a} \in \Pi_{n, k}} q^{\operatorname{maj}\left(\mathbf{a}^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[n]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q}
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From a theorem of Schensted it follows that

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From a theorem of Schensted it follows that
$\sum_{\mathbf{a} \in \boldsymbol{\Pi}_{\mathbf{n}, \mathbf{k}}} \mathbf{q}^{\operatorname{maj}\left(\mathbf{a}^{-1}\right)}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu} \sum_{\mathbf{T} \in \mathbf{S T}(\mathbf{n}-\mathbf{k}, \mu)} \mathbf{q}^{\operatorname{maj}(\mathbf{T})}$

## Do you think that was enough? no! here is more

$$
\sum_{a \in \Pi_{n, k}} q^{\operatorname{maj}\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[\mathbf{n}]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q}
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From a theorem of Schensted it follows that
$\sum_{\mathbf{a} \in \boldsymbol{\Pi}_{\mathbf{n}, \mathbf{k}}} \mathbf{q}^{\mathbf{m a j}\left(\mathbf{a}^{-\mathbf{1}}\right)}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu} \sum_{\mathbf{T} \in \mathbf{S T}(\mathbf{n}-\mathbf{k}, \mu)} \mathbf{q}^{\mathbf{m a j}(\mathbf{T})}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{S}_{\mathbf{n}-\mathbf{k}, \mu}\right\rangle$

## Do you think that was enough? no! here is more

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From a theorem of Schensted it follows that

$$
\begin{aligned}
\sum_{\mathbf{a} \in \boldsymbol{\Pi}_{\mathbf{n}, \mathbf{k}}} \mathbf{q}^{\mathbf{m a j}\left(\mathbf{a}^{-1}\right)} & =\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu} \sum_{\mathbf{T} \in \mathbf{S T}(\mathbf{n}-\mathbf{k}, \mu)} \mathbf{q}^{\mathbf{m a j}(\mathbf{T})}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{S}_{\mathbf{n}-\mathbf{k}, \mu}\right\rangle \\
& =\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{S}_{\mu}\right\rangle
\end{aligned}
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& =\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{S}_{\mu}\right\rangle \quad \text { with } \quad \mathbf{V}_{\mathbf{m}}=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}} \mathbf{h}_{\mathbf{m}+\mathbf{s}} \mathbf{e}_{\mathbf{s}}^{\perp}
\end{aligned}
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\sum_{a \in \Pi_{n, k}} q^{\operatorname{maj}\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[\mathbf{n}]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q}
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& =\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{S}_{\mu}\right\rangle \quad \text { with } \mathbf{V}_{\mathbf{m}}=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}} \mathbf{h}_{\mathbf{m}+\mathbf{s}} \mathbf{e}_{\mathbf{s}}^{\perp} \\
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\end{aligned}
$$

## Do you think that was enough? no! here is more

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\sum_{a \in \Pi_{n, k}} q^{\operatorname{maj}\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[\mathbf{n}]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q}
$$

## Proof (sketch)

the Frobenius Characteristic of Sn Harmonics

From a theorem of Schensted it follows that

$$
\begin{aligned}
\sum_{\mathbf{a} \in \boldsymbol{\Pi}_{\mathbf{n}, \mathbf{k}}} \mathbf{q}^{\operatorname{maj}\left(\mathbf{a}^{-1}\right)} & =\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu} \sum_{\mathbf{T} \in \mathbf{S T}(\mathbf{n}-\mathbf{k}, \mu)} \mathbf{q}^{\operatorname{maj}(\mathbf{T})}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{S}_{\mathbf{n}-\mathbf{k}, \mu}\right\rangle \\
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& =\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}}\right\rangle=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}}\binom{\mathbf{k}}{\mathbf{s}}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{h}_{\mathbf{n}-\mathbf{k}+\mathbf{s}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}-\mathbf{s}}\right\rangle
\end{aligned}
$$

## Do you think that was enough? no! here is more

$$
\begin{equation*}
\sum_{a \in \Pi_{n, k}} q^{\operatorname{maj}\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[\mathbf{n}]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q} \tag{*}
\end{equation*}
$$

## Proof (sketch)

the Frobenius Characteristic of Sn Harmonics
From a theorem of Schensted it follows that

$$
\begin{aligned}
& \sum_{\mathbf{a} \in \boldsymbol{\Pi}_{\mathbf{n}, \mathbf{k}}} \mathbf{q}^{\mathbf{m a j}\left(\mathbf{a}^{-1}\right)}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu} \sum_{\mathbf{T} \in \mathbf{S T}(\mathbf{n}-\mathbf{k}, \mu)} \mathbf{q}^{\operatorname{maj}(\mathbf{T})}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\underset{\mathbf{H}_{\mathbf{n}}[\mathbf{X}]}{ }, \mathbf{S}_{\mathbf{n}-\mathbf{k}, \mu}\right\rangle \\
&=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{S}_{\mu}\right\rangle \quad \text { with } \mathbf{V}_{\mathbf{m}}=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}} \mathbf{h}_{\mathbf{m}+\mathbf{s}} \mathbf{e}_{\mathbf{s}}^{\perp} \\
&=\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}}\right\rangle=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}}\binom{\mathbf{k}}{\mathbf{s}}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{h}_{\mathbf{n}-\mathbf{k}+\mathbf{s}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}-\mathbf{s}}\right\rangle \\
& \quad \text { and (*) follows easily from this }
\end{aligned}
$$

## Do you think that was enough? no! here is more

$$
\begin{equation*}
\sum_{a \in \Pi_{n, k}} q^{\operatorname{maj}\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[\mathbf{n}]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q} \tag{*}
\end{equation*}
$$

Proof (sketch)
the Frobenius Characteristic of Sn Harmonics
From a theorem of Schensted it follows that

$$
\begin{aligned}
& \sum_{\mathbf{a} \in \boldsymbol{\Pi}_{\mathbf{n}, \mathbf{k}}} \mathbf{q}^{\mathbf{m a j}\left(\mathbf{a}^{-1}\right)}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu} \sum_{\mathbf{T} \in \mathbf{S T}(\mathbf{n}-\mathbf{k}, \mu)} \mathbf{q}^{\mathbf{m a j}(\mathbf{T})}=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\underset{\mathbf{H}_{\mathbf{n}}[\mathbf{X}]}{ }, \mathbf{S}_{\mathbf{n}-\mathbf{k}, \mu}\right\rangle \\
&=\sum_{\mu \vdash \mathbf{k}} \mathbf{f}_{\mu}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{S}_{\mu}\right\rangle \quad \text { with } \quad \mathbf{V}_{\mathbf{m}}=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}} \mathbf{h}_{\mathbf{m}+\mathbf{s}} \mathbf{e}_{\mathbf{s}}^{\perp} \\
&=\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}}\right\rangle=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}}\binom{\mathbf{k}}{\mathbf{s}}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{h}_{\mathbf{n}-\mathbf{k}+\mathbf{s}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}-\mathbf{s}}\right\rangle \\
& \quad \text { and (*) follows easily from this }
\end{aligned}
$$

Problem: Get a Combinatorial interpretation when $\mathrm{Hn}[\mathrm{X}]$ is replaced by one of our Macdonald polynomials

## Do you think that was enough? no! here is more

$$
\begin{equation*}
\sum_{a \in \Pi_{n, k}} q^{\operatorname{maj}\left(a^{-1}\right)}=\sum_{r=0}^{k}(-1)^{k-r}\binom{k}{r}[\mathbf{n}]_{\mathbf{q}}[n-1]_{\mathbf{q}} \cdots[n-r+1] \mathbf{q} \tag{*}
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&=\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{V}_{\mathbf{n}-\mathbf{k}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}}\right\rangle=\sum_{\mathbf{s} \geq \mathbf{0}}(-\mathbf{1})^{\mathbf{s}}\binom{\mathbf{k}}{\mathbf{s}}\left\langle\mathbf{H}_{\mathbf{n}}[\mathbf{X}], \mathbf{h}_{\mathbf{n}-\mathbf{k}+\mathbf{s}} \mathbf{e}_{\mathbf{1}}^{\mathbf{k}-\mathbf{s}}\right\rangle \\
& \quad \text { and (*) follows easily from this }
\end{aligned}
$$

Problem: Get a Combinatorial interpretation when $\mathrm{Hn}[\mathrm{X}]$ is replaced by one of our Macdonald polynomials

## THE END

More miracles

More miracles


More miracles


More miracles


More miracles


More miracles


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More miracles


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More miracles


More miracles


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More miracles


## More miracles



## More miracles



## More miracles



## More miracles



$$
\operatorname{dim} \mathbf{M}_{\alpha_{\mathbf{1}}} \wedge \mathbf{M}_{\alpha_{2}} \wedge \cdots \wedge \mathbf{M}_{\alpha_{\mathbf{k}}}=\frac{\mathrm{n}!}{\mathbf{k}}
$$

Etc Etc Etc

## More miracles



$$
\operatorname{dim} \mathbf{M}_{\alpha_{\mathbf{1}}} \wedge \mathbf{M}_{\alpha_{2}} \wedge \cdots \wedge \mathbf{M}_{\alpha_{\mathbf{k}}}=\frac{\mathrm{n}!}{\mathrm{k}}
$$

Etc Etc Etc

## More miracles



$$
\operatorname{dim} \mathrm{M}_{\alpha_{\mathbf{1}}} \wedge \mathrm{M}_{\alpha_{\mathbf{2}}} \wedge \cdots \wedge \mathrm{M}_{\alpha_{\mathbf{k}}}=\frac{\mathrm{n}!}{\mathrm{k}}
$$

Etc Etc Etc

